# APPLICATION FOR UNITED STATES LETTERS PATENT

by

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for a

## POLYMER ACTUATOR HAVING A CIRCULAR UNIT CELL

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# POLYMER ACTUATOR HAVING A CIRCULAR UNIT CELL

## BACKGROUND

Field of the Invention

[0002] The present invention relates to microactuators having macroscopic force and displacement. More particularly the present invention relates to microactuators having a novel unit cell shape and that are formed from plastic sheets that can be stacked or laminated for stronger actuation.

Background of the Invention

[0003] Early microelectromechanical system (MEMS) microactuator arrays were typically fabricated in silicon. Despite the many favorable attributes of silicon, however, it is not always a suitable or ideal material for every application of MEMS. For instance, silicon is brittle and subject to breaking, particularly as the total device size increases. This brittleness limits devices, especially actuators, to relatively small sizes capable of only small displacements and forces. The shapes that can be realized in silicon are typically restricted by crystalline planes or 2-D fabrication processes, and more complicated structures often result in prohibitively high cost and low yield.

[0004] In an effort to overcome the deficiencies of silicon-based MEMS, microactuators have recently been developed using polymer sheets with integral electrodes that can be energized, causing electrostatic force to be generated and thereby causing the sheets to be attracted to one another. U.S. Patent 6,255,758 to Cabuz et al. describes a polymer microactuator array that includes a plurality of sheets that are layered upon one another.

In the embodiment described, every other sheet is flat and sheets disposed between the flat sheets are corrugated resulting in what can be considered a plurality of "linear" unit cells. Large forces are achieved by reducing the unit cell size. Unfortunately, this reduces the displacement as well. To achieve large displacements, many layers of unit cells have to be assembled on top of each other. Of course, with increased numbers of layers comes increased fabrication cost, yield reduction, and increased difficulty of assembly and electrical interconnection. Also, only about 1/3 of the input electrical energy does useful work. While a polymer actuator structure like that described in Cabuz et al. is demonstrably "better" than the conventional silicon based devices, there is still a desire to further improve polymer actuators.

## BRIEF SUMMARY OF THE INVENTION

[0005] The present invention provides a significant improvement over conventional "linear" type cells employed in polymer actuators or MEMS. As is described in detail later herein, by changing the unit cell shape from linear to circular significant performance improvements can be achieved. In particular, a circular unit cell can be designed to virtually any force specification and independently to virtually any displacement specification. The circular unit cell also has a theoretical efficiency of nearly 100%, so that nearly all electrical input energy is converted to useful work.

[0006] In accordance with a preferred embodiment of the present invention there is provided a microactuator device comprising at least a pair of polymeric sheets each having conductive and dielectric films deposited thereon, the polymeric sheets facing each other and bonded together to create at least one cell having a substantially circular shape substantially parallel to a plane in which the polymeric sheets lie, the at least one

2

cell having at least one egress hole to allow a fluid to pass there through when a source of electric potential is applied to the conductive films to cause a portion of the polymeric sheets in the vicinity of a perimeter of the cell to be attracted to one another and thereby cause the cell to retract upon itself.

[0007] The features and attendant advantages of the present invention will become more apparent to those skilled in the art upon a reading of the following detailed discussion in conjunction with the associated drawings.

## BRIEF DESCRIPTION OF THE DRAWINGS

[0008] Figure 1 is a plan view of circular unit cells in a polymer actuator in accordance with the present invention.

[0009] Figure 2 is a cross sectional view of one of the circular cells in accordance with the present invention.

[0010] Figures 3A and 3B depict two possible configurations for a circular cell in accordance with the present invention.

[0011] Figure 4 is a graph showing force versus displacement for a linear unit cell from equations 5 and 10.

[0012] Figure 5 is a graph showing displacement versus voltage for a linear unit cell.

[0013] Figure 6A is a graph comparing an analytical model with a finite element model (FEM) of a linear unit cell.

[0014] Figure 6B is a graph comparing displacement versus voltage calculated by the analytical model and FEM of a linear unit cell.

[0015] Figure 7 is a graph showing force versus displacement for a circular unit cell from equations 12 and 17.

3

## DETAILED DESCRIPTION OF THE INVENTION

[0016] Figure 1 is a plan view of a polymer actuator in accordance with the present invention. Reference numeral 100 is a top sheet of polymeric material such as KAPTON®, MYLAR® or KALADEX®, like those mentioned in U.S. Patent 6,255,758, which is incorporated herein in its entirety. A bottom sheet 101 (shown in Figure 2) is preferably made of the same material selected for top sheet 100. In contrast to the linear unit cells of the prior art, the present invention provides circular unit cells 102. While Figure 1 depicts only four cells, any number of circular unit cells 102 can be provided to accommodate the particular application for which the actuator is employed.

[0017] A circular cell is provided by "sweeping" the cross section of a conventional linear unit cell into a circular shape. Consequently, a circular unit cell in accordance with the present invention is fabricated by fusing or adhesively bonding top and bottom sheets 100, 101 around a circumference or perimeter of the cell. In a preferred embodiment, at least one leak (or egress) hole 104 is provided on at least one of the top and bottom sheets 100, 101 to allow air (or another fluid) to flow from inside the cell when power is applied to electrodes causing the cell to retract onto itself. Instead of leak hole 104, a relatively small portion of the perimeter of circular unit cell 102 can be left unfused or unglued to provide the same air (or fluid) egress functionality.

[0018] Figure 2 is a cross sectional view of one of the circular cells 102 in accordance with the present invention and shows the top and bottom sheets 100, 101, adhesive 110 at a perimeter portion of the cell, a metal electrode 112 and a dielectric 114. This type of layer structure is well-known and is described in, for example, U.S. Patent 6,255,758. A more detailed description of the present invention follows with reference to Figures 3-7.

## Analytical Modeling of a Linear Unit Cell

[0019] An analytical model of a linear unit cell is derived first. This is then modified to explore the effects of a circular unit cell shape.

[0020] Figures 3A and 3B show the parameters used in the analytical model of a linear unit cell. Two basic designs – one with a rigid (flat) centerplane and the other with two bowed sheets – are shown. Only half of the unit cell is shown. The other half is identical, by symmetry. The distance from the center of the cell to the edge is $l$. When no voltage is applied, the surfaces are separated at all points except the end point at $l$. The drawings show the cell with a voltage applied, where the contact point has moved to a distance $c$ from the center of the cell. Obviously, $c$ can never be larger than $l$.

[0021] The model uses energy minimization to derive equations for deflection with and without applied voltage, and an equation for the pull-in voltage. The "system," in this case, is the actuator plus its power source. The total energy of the system has four terms:

1. electrostatic energy, $U_E$, stored in the electric field of the capacitor formed by the metal films;
2. electrical work, $U_P$, done by the power source;
3. bending energy, $U_B$, of the actuator (i.e., the spring energy of the polymer structure);
4. work, $W=F_{ext}\delta$, done on the actuator by the external force, $F_{ext}$, moving the actuator through a distance $\delta$.

[0022] The electrostatic field (and therefore the electrostatic energy) drops off very rapidly as the separation between the sheets increases. Simple estimates and finite element models show that the field is negligible when the gap is a little over 1 μm. This occurs over a distance along the x-axis on the order of 10 microns or less, which is small compared to the total length of the cell. Therefore, we can assume, with little loss in accuracy, that the electrostatic energy is zero everywhere the gap is nonzero, and is

5

constant where the gap is zero. The gap is nonzero from the left hand edge (in Figures 3A and 3B) to the contact point $c$, and is zero beyond that. The electrostatic energy, $U_E$, is then simply that of the parallel plate capacitor in the contacted region. The work done by the power source is $U_P=-QV$ where $Q$ is the amount of charge that has had to move through the potential $V$. The sign is negative because the potential energy of the source decreases as it charges the capacitor. Since the actuator is a capacitor, $Q=CV$, and $U_P=-CV^2$. Therefore,

$$U_E + U_P = \frac{1}{2}CV^2 - CV^2 = -\frac{\varepsilon_0 \varepsilon b (l-c)V^2}{4d} \tag{1}$$

[0023] The bending energy is derived from the curvature of the sheet. Since the electrostatic force falls off so quickly, it has negligible influence on the profile of the curved sheet. This profile is that of a beam with a load $F_{ext}$ at the end, and can be found in standard texts and reference books on stress and strain in materials. For example, Roark and Young, *Formulas for Stress and Strain*, gives the profile as

$y = \frac{F_{ext}}{Ebt^3}l^3 - \frac{3F_{ext}c}{Ebt^3}x^2 + \frac{2F_{ext}}{Ebt^3}x^3$. The bending energy is given by

$$U_B = \frac{Ebt^3}{24}\int\left(\frac{\partial^2 y}{\partial x^2}\right)^2 dx = \begin{cases} \frac{1}{2}k\left(\frac{l}{c}\right)^3\delta^2 & c < l \\[2ex] \frac{1}{2}k\delta^2 & c = l \end{cases} \tag{2}$$

where k is the "spring constant" of the bent sheet:

$$k \equiv \frac{F_{ext}}{\delta} = \frac{Ebt^3}{nl^3} \tag{3}$$

$n=1$ if the structure has flat centerplanes, as in Figure 3A, and $n=2$ if the flat planes are missing, as in Figure 3B.

**[0024]** The total energy is then

$$U_T = \frac{1}{2} k_A \frac{l^3}{c^3} \delta^2 - \frac{\varepsilon_0 \varepsilon b(l-c)V^2}{4d} + F_{ext} \delta, \qquad (4)$$

**[0025]** First, consider the simple cases when $V=0$. Equilibrium is found by minimizing the energy with respect to $\delta$.

$$\frac{\partial U_T}{\partial \delta} = k_A \delta + F_{ext} = 0 \qquad or \qquad F_{ext} = -k_A \delta, \qquad (5)$$

the standard expression of Hooke's Law for a spring. Note that if $V=0$, then $c=l$ since there is no electrostatic force to pull the contact point in.

**[0026]** To solve for the behavior with nonzero voltage, we first need to find the value of $c$ for a given force and voltage. This is done by minimizing the energy with respect to $c$ while holding all other dimensions fixed.

$$\frac{\partial U_T}{\partial \delta} = -\frac{3}{2} k_A \frac{l^3}{c^4} \delta^2 + \frac{\varepsilon_0 \varepsilon b V^2}{4d} = 0 \qquad (6)$$

**[0027]** Solving for $c$ gives

$$c^4 = \frac{6 d k_A l^3 \delta^2}{\varepsilon_0 \varepsilon b V^2}. \qquad (7)$$

**[0028]** Plugging this back in gives the total energy, when $c$ is at equilibrium, as

$$U_T = 2(k_A l^3)^{1/4} \left( \frac{\varepsilon_0 \varepsilon b V^2}{6d} \right)^{3/4} \delta^{1/2} - \frac{\varepsilon_0 \varepsilon b l V^2}{4d} + F_{ext} \delta. \qquad (8)$$

**[0029]** Finally, as in the $V=0$ case, the equilibrium displacement is found by minimizing the energy with respect to $\delta$.

7

$$\frac{\partial U_T}{\partial \delta} = \left(k_A l^3\right)^{1/4}\left(\frac{\varepsilon_0 \varepsilon b V^2}{6d}\right)^{3/4}\frac{1}{\delta^{1/2}} - \frac{\varepsilon_0 \varepsilon b l V^2}{4d} + F_{ext} = 0 \qquad (9)$$

which rearranges to give an expression relating the external force, the displacement,

and the applied voltage.

$$F_{ext} = -\left(\frac{\varepsilon_0 \varepsilon b k_A l}{6d}\right)^{3/4}\frac{V^{3/2}}{\left(k_A \delta\right)^{1/2}}. \qquad (10)$$

[0030] The right hand side of (10) is the combined electromechanical force, $F_{em}$,

produced by the actuator at displacement $d$ and voltage $V$. The negative sign indicates

the direction of the force (i.e., opposite to the direction of the external force). The sign

can be dropped if we are just considering the magnitude of the force. Note that this

expression is valid only for $c<l$. If $c=l$, there is no surface length in contact, and the

electrostatic energy term is identically zero. Therefore, for $c=l$ the purely mechanical

equation (5) is the correct expression.

[0031] To see that this results in a pull-in behavior, compare the two lines in Figure 4.

Initially, with no external force applied and no voltage, the actuator has zero

displacement. For illustration, suppose a 1 mN force is applied (arrow #1). With the

voltage still off, the actuator behaves like a spring, and displacement travels up the line

(arrow #2) until it reaches arrow #3, the displacement at 1 mN force (~485 μm). The

curves represent the combined electromechanical behavior when various voltages are

applied. The electromechanical force generated by the actuator at 10 V (the bottom

curve) at 485 μm is small compared to the external force and the spring-like (i.e.,

mechanical only) actuator force. Therefore, the actuator still behaves like a spring, and

the displacement does not move from position 3. The same is true at 20 V, 30 V, and 40

V. When the voltage reaches 47 V, the electromechanical curve intersects the

8

mechanical curve at 485 μm, causing the displacement to begin to get smaller (along

arrow #4). As it does, the electromechanical force increases even more, pulling the

displacement in faster and faster until the actuator is fully closed. The deflection as a

function of voltage, showing the classic pull-in behavior, is plotted in Figure 5. The

deflection is approximately constant, and $c=l$, for $V$ less than some threshold, or pull-in,

voltage.

[0032] The pull-in voltage, $V_{PI}$, is therefore the voltage at which equation (5), the purely

mechanical response, equals equation (10), the combined electromechanical response.

This yields

$$V_{PI} = \left( \frac{6d}{\varepsilon_0 \varepsilon b k_A l} \right)^{1/2} F_{ext} \tag{11}$$

Comparison with FEM

[0033] Figure 6A is a comparison of the analytical model with the finite element model.

This graph should be interpreted the same as Figure 4. The dots represent FEM results

while the lines are from the analytical model. The excellent agreement is obvious.

Figure 6B compares displacement versus voltage calculated by the analytical model and

FEM, also showing a good match. The analytical model tends to overestimate the

displacement and the pull-in voltage by a small amount, typically no more than 10-15%.

Analytical Modeling of a Circular Unit Cell

[0034] In accordance with the present invention, the unit cell can be a circle, or even a

square, with the sheets bonded together all around the perimeter. Leak holes are

preferably placed in this type of cell to allow air to flow in and out as the cell opens and

closes. Since there is more length with zero gap, this cell is believed to be stronger. The

9

model is derived the same as before, with the bending and electrostatic terms modified to account for the new geometry. Figure 3B can still be used to describe the cell, but now the profile is swept into a circle of radius *l* rather than extruded into the page. A small length, $r_o$, is defined as an area over which the external force is applied. The deflection profile for a circular plate unit cell (two individual plates) is given by (again from the Roark and Young reference)

$$\delta = 2z = -\frac{3F_{ext}(1-v^2)}{2\pi Et^3}\left[l^2 - r^2\left(1 + 2\ln\left(\frac{l}{r}\right)\right)\right] \tag{12}$$

[0035] The bending energy for a single plate is

$$dU_B = \frac{Et^3}{24(1-v^2)}\left[\left(\frac{\partial^2 z}{\partial r^2} + \frac{1}{r}\frac{\partial z}{\partial r}\right)^2 + 2(1-v)\left(\frac{1}{r}\frac{\partial z}{\partial r}\frac{\partial^2 z}{\partial r^2}\right)\right]2\pi r dr \tag{13}$$

[0036] Calculating the derivatives from Eq. (12), plugging into Eq. (13) and integrating from $r_o$ to *l* yields the total bending energy:

$$U_B = \frac{\pi}{3}\frac{Et^3}{(1-v^2)}\frac{\left\{c^2 - r_o^2 - (6-2v)r_o^2\left(\ln\frac{r_o}{c}\right)^2\right\}}{\left\{c^2 - r^2\left(1 + 2\ln\frac{c}{r}\right)\right\}^2}\delta^2$$

[0037] Assuming $r_o \ll c$, and replacing *l* with the variable parameter *c*, this simplifies to a more manageable form

$$U_B \approx \frac{\pi}{3}\frac{Et^3}{(1-v^2)c^2}\delta^2 \tag{14}$$

[0038] The electrostatic energy is based on the same assumptions as the basic model. For the circular geometry, this becomes

$$U_E + U_P = -\frac{\varepsilon_o\varepsilon\pi(l^2 - c^2)}{4d}V^2 \tag{15}$$

so that

10

$$U_T = \frac{\pi}{3} \frac{Et^3}{(1-v^2)c^2} \delta^2 - \frac{\varepsilon_o \varepsilon \pi (l^2 - c^2)}{4d} V^2 + F_{ext} \delta \,. \tag{16}$$

[0039] Following the same procedure as before, we take the derivative with respect to $c$, solve for $c$, plug that value back into the total energy, and finally set $\partial U_T / \partial \delta = 0$, with the result

$$F_{ext} = -\frac{\pi\sqrt{3}}{4} \left( \frac{Et^3}{(1-v^2)} \frac{\varepsilon_o \varepsilon}{4d} \right)^{1/2} V \,. \tag{17}$$

[0040] This is analogous to the linear unit cell model equation (10), and is valid for $c<l$. The interesting feature of equation (17) is that there is no dependence on $\delta$. Figure 7 shows the force-displacement graph for the circular cell, analogous to Figure 4 for the linear cell. When the external force is applied, the cell opens up to some displacement. As the voltage is turned on and increased a combined electromechanical response is generated which is flat as a function of $\delta$. Initially, the electromechanical force is less than the external force, and nothing happens. However, when the voltage is large enough, the actuator pulls in. The circular unit cell pull-in is different than the linear unit cell pull-in, however. In the model of the linear unit cell, as the cell pulled in slightly, the electromechanical force got even larger, producing a runaway pull-in. In the circular unit cell, the electromechanical and external forces are equal, resulting in no pull-in until the voltage just exceeds $V_{PI}$. Pull-in occurs, but it is not a runaway pull-in. Another way to express this is that in the linear model, kinetic energy is produced and increases as the cell pulls in. In the circular model, the cell pulls in at a constant velocity, and there is no increase in kinetic energy. This phenomenon has consequences when we consider power consumption later herein. The pull-in voltage comes from rearranging eq. (17)

11

$$V_{PI} = \frac{4}{\pi\sqrt{3}} \left( \frac{(1-v^2)}{Et^3} \frac{4d}{\varepsilon_o\varepsilon} \right)^{1/2} F_{ext}.$$ (18)

[0041] There are two very important consequences with respect to a circular unit cell. The first is that the force generated by a single cell of the actuator is independent of the displacement. This is very different from the linear unit cell where force varied as $\delta^{1/2}$ (see eq. 10). To generate a large force with the linear cell, one has to settle for a reduced displacement. To get large displacements, many layers of unit cells must be stacked on top of each other. This requires fabrication of many cells, some form of assembly procedure, and a means of electrically connecting to all of those layers. In contrast, the circular cell can be designed to achieve substantially any force and substantially any displacement. Force is controlled most easily by controlling the thickness (eq. 17 and 18) and displacement can be independently controlled by controlling the unit cell diameter (eq. 12). Thus, a single layer of circular unit cells (or at most a small number of layers) can be designed to do the job of many layers of linear unit cells. This greatly reduces the complexity and cost of fabrication and assembly, and improves yield.

[0042] The second consequence is that the power consumption of a circular unit cell is a minimum. The argument for this comes from the fact that the force generated by the cell is constant as it pulls in. No kinetic energy is created and, therefore, no energy is lost when motion stops at the end of the pull-in. Thus, all of the electrical energy goes in to useful work and none is wasted. Theoretically, the efficiency of this cell is nearly 100%. In reality, of course, there are losses, etc. that reduce the efficiency, but no unit cell design can be more efficient. In addition to the high efficiency, the design flexibility described in the previous paragraph also contributes to a significant energy savings. As described above, a circular-cell actuator can be constructed from a smaller number of

12

layers than a linear-cell actuator. Since the total capacitance of the actuator is approximately proportional to the number of layers, the circular-cell actuator has a lower total capacitance. This reduces the energy required to charge the capacitor and the energy lost when discharging the capacitor.

[0043] The foregoing disclosure of the preferred embodiments of the present invention has been presented for purposes of illustration and description. It is not intended to be exhaustive or to limit the invention to the precise forms disclosed. Many variations and modifications of the embodiments described herein will be apparent to one of ordinary skill in the art in light of the above disclosure. For example, the circular cells described herein need not be precisely circular, but are preferably substantially circular to achieve the performance described. The scope of the invention is to be defined only by the claims appended hereto, and by their equivalents.

[0044] Further, in describing representative embodiments of the present invention, the specification may have presented the method and/or process of the present invention as a particular sequence of steps. However, to the extent that the method or process does not rely on the particular order of steps set forth herein, the method or process should not be limited to the particular sequence of steps described. As one of ordinary skill in the art would appreciate, other sequences of steps may be possible. Therefore, the particular order of the steps set forth in the specification should not be construed as limitations on the claims. In addition, the claims directed to the method and/or process of the present invention should not be limited to the performance of their steps in the order written, and one skilled in the art can readily appreciate that the sequences may be varied and still remain within the spirit and scope of the present invention.